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Excess volumes, partial molar and adiabatic compressibilities of binary mixtures of *n*-alcohols with monoethanolamine

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Speed of sound and densities of solution mixtures of four aliphatic alcohols with monoethanolamine were measured over a full range of composition. The density measurements were carried out from 298 to 358 K. Results of these measurements were used to calculate adiabatic compressibilities, excess adiabatic compressibilities, excess volumes and partial molar quantities. From the analysis of the results, the nature of interaction between the aliphatic alcohols and monoethanolamine is deduced.

Keywords: aliphatic alcohols; speed of sound; volumes; adiabatic compressibility and partial molar properties

1. Introduction

Aqueous solutions of monoethanolamine (MEA) have been extensively used for the removal of acidic gases in the natural gas. The removal of carbon dioxide by using chemical absorbents has been of considerable interest due to excessive discharge of carbon dioxide and methane which are primarily responsible for global warming. Aqueous monoethanolamine solutions are known to react with materials in reactor vessels, tubing lines and other process compartments. Blended mixtures of MEA with other solvents may have the potential to minimise the adverse effects to reactor vessels and tubing lines. Several investigations on the removal of acid gases by blended MEA mixtures have been reported in the literature [1–4]. Recently, Lee *et al.* [5–6] reported densities and viscosities of binary solutions of MEA with aliphatic alcohols at 303, 313 and 323 K.

Knowledge of physical properties of blended alkanolamine mixtures is useful in connection with industrial treatment of acidic gases. The properties of these mixtures at several temperatures are required for engineering design of acid gas treatment equipment and for subsequent operations as summarised by Astarita *et al.* [7]. The physical properties are also needed for predicting the absorption rates and enhancement factors of the gases, which are useful to process engineers who are interested in the problem of acidic gas removal. Volumetric data in combination with molecular theories of solutions can be used to extend our understanding of molecular interactions occurring in these systems. Derived volumetric properties can be employed to arrive at a more detailed picture of solute–solvent interactions [8].

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	$ ho ({\rm gcm^{-3}}) 298.15{\rm K}$		$U(m s^{-1})$		
Compound	Expt.	Lit.	Expt. (303 K)	Lit. (298 K)	
1-Propanol Isopropanol	0.79981 0.78125	0.79975 [9] 0.78090 [10]	1199.0 1126.6	1206 [9] 1140 [10] 1127 [12]*	
1-Hexanol 1-Octanol	0.81530 0.82174	0.81534 [11] 0.82230 [12]	1288.1 1332.2	1303 [11] 1348 [12]	

Table 1. Experimental densities and speed of sound of pure compounds along with the literature values.

Note: *The value corresponds to T = 303.15 K.

The present study reports density measurements of the binary mixtures of MEA with aliphatic alcohols such as 1-propanol, isopropanol, 1-hexanol and 1-octanol in the temperature of 298–358 K along with speed of sound measurements in the entire composition range at 303 K. Various derived properties such as excess volumes ($V^{\rm E}$), thermal expansivity (α), adiabatic compressibility (β), molar adiabatic compressibility (βv) and apparent molar compressibility K_{ϕ} , are evaluated and discussed in terms of molecular interactions in the mixtures.

2. Experimental section

The ethanolamine and all aliphatic alcohols were obtained from Aldrich of stated purity at 99% by mass and used as received. The densities and sound velocity values of the pure components agreed well with the values reported in the literature [9–13]. The values are presented in Table 1. Density of mixtures were measured with an Anton Parr DMA 5000 Density meter which was calibrated with dry air and distilled water. Accuracies of our density measurements are about $\pm 3 \times 10^{-5} \text{ g cm}^{-3}$. The accuracy at higher temperatures may range around $\pm 8 \times 10^{-5} \text{ g cm}^{-3}$. Temperature was stable to $\pm 0.001^{\circ}$ C during the course of measurements. The speed of sound measurements were carried out as described in our previous publications [14–16].

3. Results and discussion

The experimental densities and speed of sound measurements of the binary mixtures over the entire composition range are presented in Tables 2 and 3, respectively.

3.1. Speed of sound

The adiabatic compressibilities were calculated using the formula

$$\beta = \frac{1}{U^2 \rho},\tag{1}$$

where U is the ultrasonic speed of sound and ρ is the density of the solution. The adiabatic compressibility values are found to be dependent on composition

Table 2. Densities $(g cm^{-3})$ of binary mixture at various temperatures.

X _{MEA}	298 K	308 K	318 K	328 K	338 K	348 K	358 K
Propanol and MEA							
0.00000	0.79981	0.79169	0.78342	0.77492	0.76617	0.75709	0.74762
0.05028	0.80932	0.80120	0.79293	0.78446	0.77575	0.76675	0.75739
0.09906	0.81877	0.81064	0.80237	0.79393	0.78526	0.77632	0.76706
0.20023	0.83232	0.82332	0.81390	0.80382	0.79297	0.78015	0.76925
0.29995	0.85761	0.84783	0.83668	0.82413	0.80955	0.79048	0.78302
0.40102	0.87964	0.87151	0.86330	0.85497	0.84649	0.83778	0.81744
0.49961	0.90049	0.89239	0.88420	0.87590	0.86748	0.85891	0.85016
0.60024	0 92214	0 91405	0.90588	0.89763	0.88927	0.88077	0.87212
0.70078	0.94446	0.93639	0.92826	0.92005	0.91175	0.90333	0.89477
0 79984	0.96682	0.95878	0.95069	0.94252	0.93428	0.92593	0.91745
0.89890	0.98922	0.98121	0.97316	0.96505	0.95687	0.94860	0.94021
0.07062	1 00051	0.00252	0.08450	0.90505	0.96827	0.94000	0.05170
1,00000	1.00051	1.00503	0.98430	0.97042	0.90827	0.90003	0.95170
1.00000	1.01299	1.00505	0.99704	0.98899	0.98088	0.97209	0.90441
Isopropanol and MEA							
0.00000	0.78125	0.77261	0.76362	0.75423	0.74434	0.73391	0.72286
0.05082	0.79134	0.78277	0.7739	0.76469	0.75503	0.74488	0.73360
0.09972	0.80109	0.79258	0.78381	0.77473	0.76525	0.75536	0.74500
0.20334	0.82521	0.81680	0.80821	0.79937	0.79025	0.78080	0.77095
0.29817	0.84374	0.83539	0.82687	0.81816	0.80921	0.79996	0.79038
0.40261	0.86701	0.85873	0.85032	0.84174	0.83296	0.82395	0.81467
0.50055	0.89020	0.88197	0.87364	0.86517	0.85653	0.84768	0.83864
0.60015	0.91362	0.90545	0.89719	0.88881	0.87912	0.86105	0.83832
0.69914	0.93787	0.92975	0.92156	0.91327	0.90486	0.89313	0.88760
0.79941	0.96233	0.95425	0.94611	0.93789	0.92959	0.92116	0.91261
0 90015	0 98778	0 97975	0 97168	0.96355	0.95534	0.94703	0.93862
0.94863	1 00010	0.99210	0.98406	0 97597	0.96781	0.95956	0.95121
1.00000	1.01299	1.00503	0.99704	0.98899	0.98088	0.97269	0.96441
Hexanol and MEA	1.012	1.00202	0.55701	0.90099	0.90000	0.97209	0.90111
0.00000	0.81530	0.80807	0.80072	0 79321	0 78553	0 77765	0 76954
0.05658	0.82077	0.81351	0.80614	0.79865	0.79098	0.78312	0.77504
0.10568	0.82589	0.81860	0.81121	0.80370	0.79604	0.78819	0.78014
0.20224	0.83692	0.82956	0.82213	0.81458	0.80689	0.70012	0.70014
0.20224	0.83072	0.82750	0.82215	0.82634	0.81861	0.81073	0.80267
0.29802	0.84087	0.85652	0.83334	0.82034	0.81301	0.81075	0.80207
0.50170	0.80403	0.83032	0.84894	0.84127	0.83349	0.82337	0.81748
0.50170	0.00014	0.87233	0.00409	0.03/10	0.04932	0.84133	0.85523
0.39897	0.89843	0.89078	0.88500	0.87323	0.80/30	0.83934	0.85120
0.70032	0.92059	0.91284	0.90505	0.89/18	0.88923	0.88117	0.8/299
0.20014	0.945/5	0.93/94	0.93007	0.92214	0.91413	0.90603	0.89/81
0.89814	0.97572	0.96/83	0.95991	0.95192	0.94387	0.93572	0.92/4/
0.94670	0.99284	0.98492	0.9/696	0.96896	0.96087	0.952/1	0.94444
1.00000	1.01299	1.00503	0.99704	0.98899	0.98088	0.97269	0.96441
Octanol and MEA							
0.00000	0.82174	0.81476	0.80769	0.80050	0.79317	0.78568	0.77802
0.05011	0.82515	0.81815	0.8110	0.80387	0.79655	0.78906	0.78141
0.10065	0.82922	0.82218	0.81507	0.80787	0.80053	0.79305	0.78539
0.20057	0.83797	0.83086	0.82368	0.81641	0.80903	0.80151	0.79384
0.30340	0.84815	0.84094	0.83367	0.82633	0.81888	0.81131	0.80358
0.40153	0.85971	0.85239	0.84503	0.83759	0.83006	0.82242	0.81463
0.49814	0.87234	0.86493	0.85747	0.84994	0.84233	0.83460	0.82673
0.59863	0.88947	0.88194	0.87437	0.86674	0.85902	0.85120	0.84326
0.70210	0.91081	0.90316	0.89547	0.88773	0.87990	0.87198	0.86395
0.80010	0.93595	0.92819	0.92039	0.91254	0.90462	0.89660	0.88849
0 90016	0.96890	0.96104	0.95314	0.94520	0.93718	0.92908	0.92089
0 94911	0.98893	0.98103	0 97310	0.96511	0 95706	0.94893	0.94070
1 00000	1 01299	1 00503	0 99704	0 98899	0.98088	0 97269	0 96441
	1.014//	1.000000	0.77701	0.,00,0	0.20000	0.77207	0.20171

	X _{ME}
	0.000
	0.099
	0.144
	0.201
	0.251
	0.305
	0.342
	0.385
	0.407
	0.448
	0.469
	0.489
	0.512
	0.533
	0.584
	0.617
-	0.652
201	0.687
rУ	0.720
nua	0.749
Jai	0.783
28	0.821
30	0.862
: 40	0.899
 L	0.933
A.	0.955
Idec	0.971
loa	0.988
Down	1.000

Table 3. Speed of sound of binary mixtures MEA and alcohols at 303 K.

Propanol		Isopropanol		He	exanol	Octanol	
X _{MEA}	$U (\mathrm{ms}^{-1})$	X _{MEA}	$U (\mathrm{ms}^{-1})$	X _{MEA}	$U (\mathrm{ms}^{-1})$	X _{MEA}	U/ms^{-1}
0.0000	1199.0	0.0000	1126.6	0.0000	1288.1	0.0000	1332.2
0.0994	1232.1	0.0111	1130.6	0.0267	1293.3	0.0307	1336.3
0.1441	1250.6	0.0301	1138.4	0.1077	1304.2	0.1305	1344.1
0.2016	1275.9	0.0476	1146.1	0.1777	1316.8	0.1868	1349.5
0.2515	1300.5	0.0893	1164.4	0.2217	1324.4	0.2360	1356.4
0.3052	1323.8	0.1047	1171.3	0.2628	1333.1	0.2811	1363.6
0.3425	1341.4	0.1392	1180.9	0.2928	1343.6	0.3022	1367.5
0.3852	1360.5	0.1787	1199.3	0.3365	1353.2	0.3448	1372.2
0.4070	1370.9	0.2031	1210.3	0.3832	1365.2	0.3833	1377.7
0.4487	1391.3	0.2289	1224.9	0.4149	1373.3	0.4108	1385.5
0.4690	1401.5	0.2487	1233.9	0.4461	1381.3	0.4419	1391.4
0.4895	1413.7	0.2667	1244.2	0.4702	1388.9	0.4715	1397.4
0.5129	1424.6	0.2957	1259.3	0.4983	1396.6	0.5090	1406.3
0.5330	1436.2	0.3379	1282.2	0.5296	1407.3	0.5491	1415.9
0.5847	1462.4	0.3551	1291.8	0.5524	1415.6	0.5863	1426.4
0.6175	1479.9	0.4066	1315.8	0.5669	1421.2	0.6314	1439.8
0.6524	1499.7	0.4440	1337.5	0.5879	1428.7	0.6628	1451.0
0.6870	1519.0	0.4666	1357.7	0.6180	1441.2	0.6938	1463.0
0.7201	1536.5	0.5126	1386.7	0.6376	1449.0	0.7176	1472.6
0.7491	1552.8	0.5609	1421.0	0.6581	1459.1	0.7444	1484.7
0.7830	1573.6	0.5874	1431.8	0.7099	1479.7	0.7754	1502.6
0.8219	1595.2	0.6102	1447.6	0.7438	1498.0	0.7895	1510.4
0.8620	1620.3	0.6583	1477.3	0.7829	1522.0	0.8206	1529.0
0.8998	1641.5	0.7143	1511.8	0.8309	1553.8	0.8489	1549.3
0.9334	1664.3	0.7826	1557.3	0.8756	1588.2	0.8820	1575.6
0.9552	1675.7	0.8140	1578.6	0.9249	1629.5	0.9223	1623.6
0.9718	1685.2	0.8563	1607.4	0.9582	1660.0	0.9532	1669.1
0.9886	1694.1	0.8938	1626.4	0.9905	1692.4	0.9762	1682.2
1.0000	1703.7	0.9097	1642.6	1.0000	1703.7	0.9922	1694.9
		0.9438	1666.2			1.0000	1703.7
		0.9759	1688.3				
		1.0000	1703.7				

as well as on the hydrophobic character of the aliphatic alcohol. Excess adiabatic compressibility coefficients as a measure of deviation from ideal behaviour are obtained as follows:

$$\Delta \beta = \beta_{\text{expt}} - \beta_{\text{ideal}}.$$
 (2)

The ideal values (β_{id}) are calculated by taking the volume-fraction ϕ_1 weighted contributions of the isothermal compressibilities [17–20]. From the relation between isothermal and adiabatic compressibilities, (β_{id}) can be expressed as follows:

$$\beta_{id} = \sum \phi_i \left(\beta_i + \frac{TV_i^* \alpha^2}{C_{P,i}} \right) - \frac{T\left(\sum \left(X_1 V_i^* \right) \sum \left(\phi_i \alpha_i \right)^2 \right)}{\sum \left(X_i C_{P,i} \right)},\tag{3}$$



Figure 1. Plot of excess isentropic compressibilities ($\Delta\beta$) as a function of MEA mole fraction for various alcohols at 303 K.

Table 4. $(\beta)^{E}/10^{-12}$ Redlich parameters for MEA + alcohol at 303.15 K.

	Propanol	Isopropanol	Hexanol	Octanol
A	-352.85	-527.00	-122.86	-35.50
В	8.70	57.31	61.40	31.74
С	-0.42	141.31	-18.80	-54.79
D	66.21	-21.62	53.09	64.84
Е	91.43	-152.36	-22.35	_
R^2	0.999	0.999	0.997	0.986

where β_i , V_i^* , α_i and $C_{P,i}$ are the adiabatic compressibility, molar volume, coefficient of thermal expansivity and molar heat capacity at a constant pressure of the *i*th component, respectively. The heat capacity data for the pure components is taken from the literature [21,22]. The excess property $\Delta\beta$ calculated from Equations (2) and (3) were fitted to Redlich–Kister equation [23]

$$y^{E} = X_{1}X_{2}\sum_{j=0}^{n} a_{j}(1 - 2X_{1})^{j},$$
(4)

where y^{E} and a_{j} represent the excess property in question and the polynomial coefficient of the fitting equation.

The $\Delta\beta$ plots at 303 K are shown in Figure 1 and the fitting coefficients of Equation (4) are presented in Table 4. The thermal expansivity values (α_i) were calculated from the dependence of molar volume as a function of temperature and is given by

$$\alpha_i = \frac{1}{V_i^*} \left(\frac{\partial V_i}{\partial T} \right)_P. \tag{5}$$

X _{MEA}	Propanol	Isopropanol	Hexanol	Octanol
0.0	4.80	-5.76	15.22	21.36
1.0	44.93	46.28	79.59	101.16

Table 5. $K_{\phi} \times 10^{15}$ for MEA + alcohol at 303.15 K.

The $\Delta\beta$ values were found to be negative in all cases and a pronounced minimum was observed around 0.5 mole fraction of the MEA for hydrophilic propanol and isopropanol. The negative trends indicate the strength of interaction between these species. In the case of more hydrophobic alcohols, only weak interactions are expected, resulting in similar negative values for $\Delta\beta$.

The excess molar adiabatic compressibility $(\beta V)^{\rm E}$ or $K^{\rm E}$ is given by

$$(\beta V)^E = \beta_s V - (\beta_s V)^{id}, \tag{6}$$

$$(\beta_s V)^{id} = X_1 \beta_1^* V_1^* + (1 - X_1) \beta_2^* V_2^*, \tag{7}$$

where β_s , β_1^* , β_2^* , V, V_1^* , and V_2^* , are adiabatic compressibilities and volumes of the solution of components 1 and 2, respectively. The molar volumes of solutions were calculated using the following equation:

$$V = \frac{M_1 X_1 + M_2 (1 - X_1)}{\rho},\tag{8}$$

where M_1 and M_2 are the molecular weights of components 1 and 2, respectively. The apparent molar compressibilities of components were obtained from $(\beta V)^E$ using the relations [24]

$$K_{\phi,1} = \frac{(\beta V)^{\rm E}}{X_1} + \beta_1^* V_1^* \tag{9}$$

and

$$K_{\phi,2} = \frac{(\beta V)^E}{X_2} + \beta_2^* V_2^*.$$
 (10)

From the plots of $K_{\phi,1}$ and $K_{\phi,2}$ as a function of mole fraction of MEA, the partial molar compressibility values at infinite dilution $K_{\phi,1}^0$ and $K_{\phi,2}^0$ were obtained by means of graphical methods and are listed in Table 5. The values of $K_{\phi,2}^0$ increase in the order of increasing hydrophobic nature of the alcohols [25,26].

3.2. Volumes

The excess volume is given by:

$$V^E = V - V^{id},\tag{11}$$

where $V^{id} = V_1^* X_1 + (1 - X_1) V_2^*$ The excess volumes were plotted as a function of mole fraction of X_{MEA} at all temperatures and presented in Figure 2. The fitting parameters of the plots using Equation (4) are presented in Table 6. The excess volumes were found to be



Figure 2. Plots of excess volumes for the systems MEA with alcohols as a function of temperature.

Parameters	298 K	308 K	318 K	328 K	338 K	348 K	358 K
Propanol							
a_0	-1.7762	-1.8272	-1.8915	-1.9784	-2.0869	-2.2166	-2.3857
a_1	0.1479	0.1593	0.1405	0.0876	0.0794	0.1470	0.1694
a_2	0.1741	0.1932	0.1746	0.1363	0.1455	0.1432	0.1347
a_3	-0.3090	-0.3174	-0.3222	-0.3001	-0.3216	-0.4740	-0.6175
R^2	0.9957	0.9958	0.9959	0.9957	0.9961	0.9971	0.9975
Isopropanol							
a_0	-2.0940	-2.2286	-2.4096	-2.6485	-2.8687	-3.2200	-3.8790
a_1	0.5453	0.5133	0.4374	0.3571	0.0355	-0.3558	-0.1041
a_2	0.2369	0.2113	0.1427	0.0946	-0.0895	-0.2254	0.1634
a_3	0.0897	0.0656	0.0407	1.0411e - 3	0.3715	0.6982	0.1642
R^2	0.9978	0.9980	0.9982	0.9984	0.9909	0.9973	0.9979
Octanol							
a_0	-5.1918	-5.2234	-5.2745	-5.2721	-5.3252	-5.6010	-5.7743
a_1	1.4604	1.4766	1.4822	1.4763	1.4465	1.4275	1.4026
a_2	-0.4819	-0.5227	-0.5739	-0.5460	-0.6233	-0.7979	-0.8802
<i>a</i> _{3_}	0.1341	0.1259	0.1072	0.0938	0.1057	9.4932e-3	-0.0420
R^2	0.9957	0.9982	0.9961	0.9973	0.9979	0.9973	0.9959

Table 6. Redlich–Kister coefficients for V^{E} of MEA and alcohol mixtures at various temperatures.

negative, indicating negative deviation from ideality with a minimum at 0.5 mole fraction for 1-propanol and isopropanol, and around 0.6 for octanol. The excess volumes calculated for the 1-hexanol system are found to be negative within the range of experimental error and therefore were not fitted to Equation (4). The excess volumes ($V^{\rm E}$) are also found to be dependant on temperature and became more negative at higher temperatures. The magnitude of $V^{\rm E}$ reflects the compactness. Negative $V^{\rm E}$ values indicate that more efficient packing is evident in the mixture compared to the pure liquid components. Thus, it can be assumed that the most hydrophobic alcohol interacts with MEA in a more compact form than the other hydrophilic alcohols.

4. Conclusions

Volumetric properties of aliphatic alcohols with MEA are found to depend on the nature of the hydrophobicity and hydrogen bonding ability of the alcohols. Alcohols with higher hydrophobic character interact more strongly with MEA, leading to larger values of $K_{\phi,2}^0$.

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References

- [1] B. Hawrylak, S.E. Burke, and R. Palepu, J. Soln. Chem. 29, 575 (2000).
- [2] Y. Maham, T.T. Teng, L.G. Hepler, and A.E. Mather, J. Soln. Chem. 23, 195 (1994).
- [3] J.H. Song, S.B. Park, J.H. Yoon, and H. Lee, J. Chem. Eng. Data 42, 143 (1997).
- [4] B.E. Roberts and A.E. Mather, Chem. Eng. Commun. 72, 201 (1998).
- [5] M.J. Lee and T.K. Lin, J. Chem. Eng. Data 40, 336 (1995).
- [6] M.J. Lee, T.K. Lin, Y.H. Pai, and K.S. Lin, J. Chem. Eng. Data 42, 854 (1997).
- [7] G. Astarita, D.W. Savage, and A. Bisio, *Gas Treating with Chemical Solvents* (Wiley, New York, 1983).
- [8] M. Sakurai, K. Nakamura, K. Nitta, and N. Takenaka, J. Chem. Eng. Data 40, 301 (1995).
- [9] E. Langa, A.M. Maihar, J.I. Pardo, and J.S. Urieta, J. Chem. Eng. Data 50, 1255 (2005).
- [10] G. Savaroglu and E. Aral, Fluid Phase Equilibr. 215, 253 (2004).
- [11] E. Mascato, L. Mosteiro, M.M. Pineiro, J. Garcia, T.P. Iglesias, and J.L. Legido, J. Chem. Thermodyn. 33, 1081 (2001).
- [12] E. Mascato, L. Mosteiro, M.M. Pineiro, J. Garcia, T.P. Iglesias, and J.L. Legido, J. Chem. Eng. Data 45, 1153 (2000).
- [13] D. Venkatesulu, P. Venkatesu, and M.V. Prabhakara Rao, J. Chem. Eng. Data 42, 1145 (1997).
- [14] B. Hawrylak, K. Gracie, and R. Palepu, Can. J. Chem. 76, 464 (1998).
- [15] S. Smith, P. Wiseman, and L. Boudreau, J. Soln. Chem. 23, 207 (1994).
- [16] K. Gracie, D. Turner, and R. Palepu, Can. J. Chem. 74, 1616 (1996).
- [17] O. Kiyohara and G.C. Benson, J. Chem. Thermodyn. 1, 861 (1979).
- [18] G. Douheret, C. Moreau, and A. Viallard, Fluid Phase Equilibr. 22, 227 (1985).
- [19] G.C. Benson and O. Kiyohara, J. Chem. Thermodyn. 11, 1061 (1979).
- [20] N.P. Rao and R.R. Verral, Can. J. Chem. 65, 810 (1987).

- [21] Y. Maham, L.G. Hepler, A.E. Mather, A.W. Hakin, and R.A. Marriot, J. Chem. Soc. Faraday Trans. **93**, 1747 (1997).
- [22] M. Bures, M. Zabransky, and V. Svoboda, Thermochim. Acta 245, 145 (1994).
- [23] O. Redlich and A.T. Kister, Ind. Eng. Chem. 40, 345 (1948).
- [24] A. Cipiciani, G. Onori, and G. Savelli, Chem. Phys. Lett. 143, 505 (1998).
- [25] S. Nishikawa and M. Mashima, J. Chem. Soc. Faraday Trans. 86, 3567 (1990).
- [26] K. Nakanishi, N. Kato, and M. Maruyama, J. Phys. Chem. 71, 814 (1967).